

Lawrence Livermore National Laboratory

Predicting properties of plutonium metal and alloys within the dynamical mean field theory

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Computing Grand Challenge Symposium



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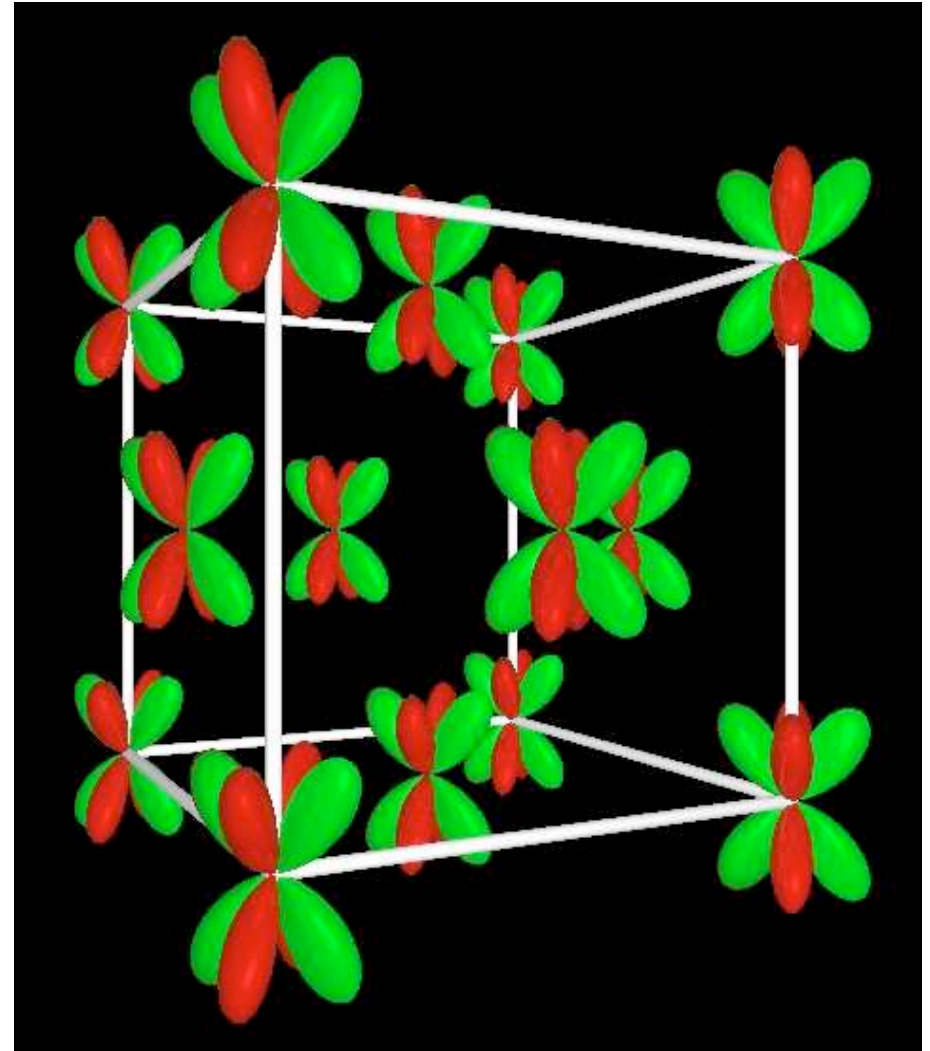
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Outline

- Computational and scientific goals
- Methodology
 - Density functional theory (DFT)
 - Dynamical mean-field theory (DMFT)
 - Continuous time QMC approach (CTQMC)
- DFT+DMFT computed properties of δ Pu
- Future work
- Conclusion



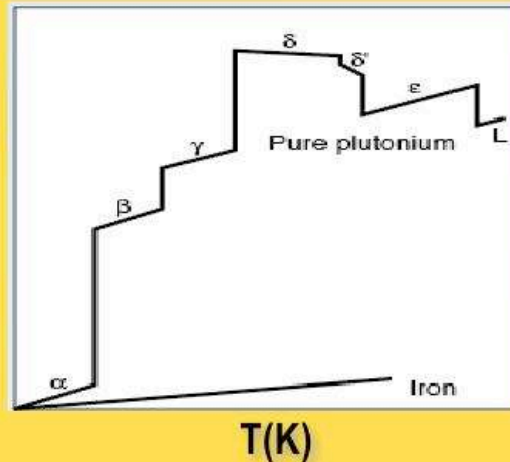
A schematic of f_{xyz} orbitals in δ Pu

Anomalous Properties of Pu

- Scientific objectives
 - Absence of magnetism
 - Volume collapse
 - Role of alloying
- Electronic properties
 - Spectra
 - Heat capacity
 - Magnetic susceptibility
 - Fermi surface
 - Resistivity

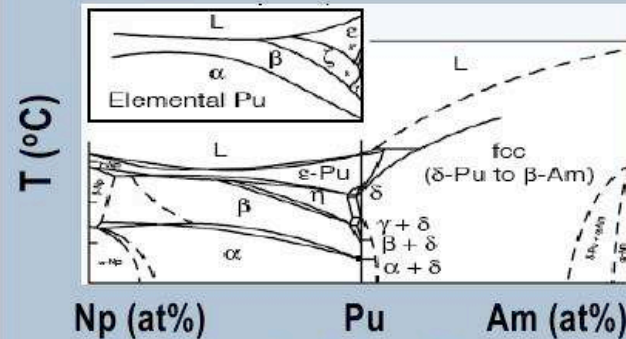
Pu and its simple alloys are metallurgical anomalies

Length change (%)



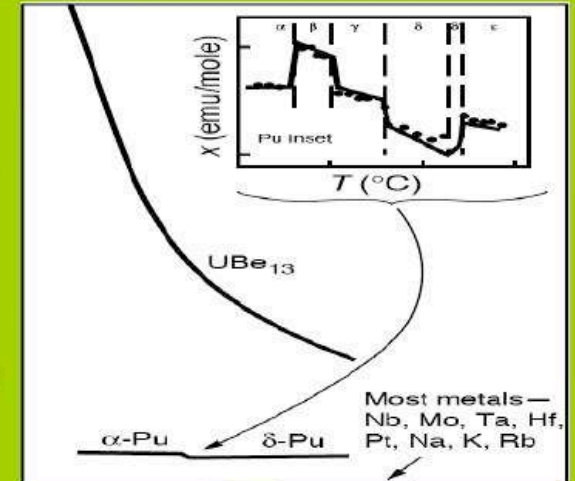
Six phases & large volume collapse

Pressure (kbar)



Complex Phase Diagrams

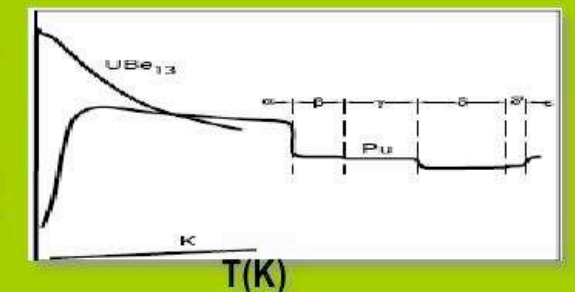
χ (emu/mole)



T(K)

Anomalous susceptibility

ρ ($\mu\Omega$ cm)



T(K)

Anomalous resistivity

DMFT is positioned to meet this challenge



General Methods

- Density functional theory (DFT)

$$\Rightarrow \Gamma[\rho]$$

- Realistic systems

- Ground state properties

- Dynamical mean-field theory (DMFT)

$$\Rightarrow \Gamma[G]$$

- Model Hamiltonians

- Captures physics of Mott transition

- Ground/excited state properties

Merge approaches \Rightarrow DFT+DMFT

$$\Rightarrow \Gamma[\rho, G]$$



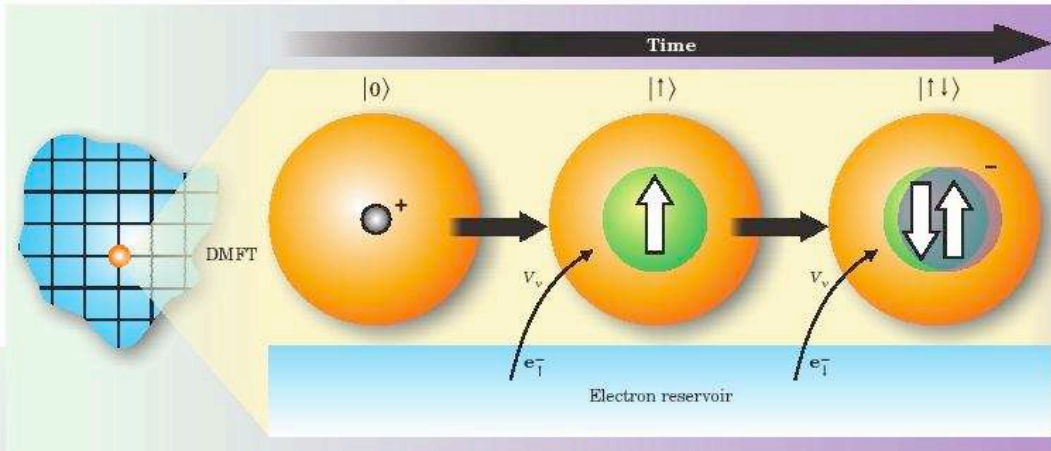
DFT+DMFT for Pu

- DMFT too expensive to apply to all electrons
 - Even for Blue Gene!
- Only apply DMFT to most correlated electrons
 - f -electrons in Pu.
 - Apply DFT to the remaining electrons.
- Double-counting problem
 - Implication: number of f -electrons not accurately predicted.
 - Comparison with experiment must determine f -electron count

Nearly first-principles approach to strongly correlated materials



DMFT: what is it?



$$H_{AIM} = \sum_{\sigma} \epsilon_{imp} c_{\sigma}^{\dagger} c_{\sigma} + U c_{\uparrow}^{\dagger} c_{\uparrow} c_{\downarrow}^{\dagger} c_{\downarrow} + \sum_{k\sigma} V_k (a_{k\sigma}^{\dagger} c_{\sigma} + c_{\sigma}^{\dagger} a_{k\sigma}) + \sum_{k\sigma} \alpha_k a_k^{\dagger} a_k$$

- *DMFT maps lattice many-body problem to Anderson impurity model (AIM).*
- AIM has 1 site of lattice embedded in bath of fictitious electrons which mimic removed lattice sites.
- Characteristics of fictitious electrons determined by DMFT self-consistency condition.
- AIM may be solved accurately using computation.



Solving Anderson impurity model

- Quantum Monte-Carlo may be used to exactly solve AIM
 - Computational cost increases as temperature decreases
- Traditional Hirsch-Fye QMC method has limitations
 - Difficult to reach ambient temperatures
 - Cannot exactly treat on-site exchange
- Continuous time QMC (CTQMC) does not have these limitations
 - CTQMC stochastically sums the Feynman diagrams of the AIM.
 - CTQMC may be applied starting from the band limit or the atomic limit

Computation can provide the exact solution to the AIM!



CTQMC: band limit vs. atomic limit

$$U = 0 \text{ vs. } V_k = 0$$

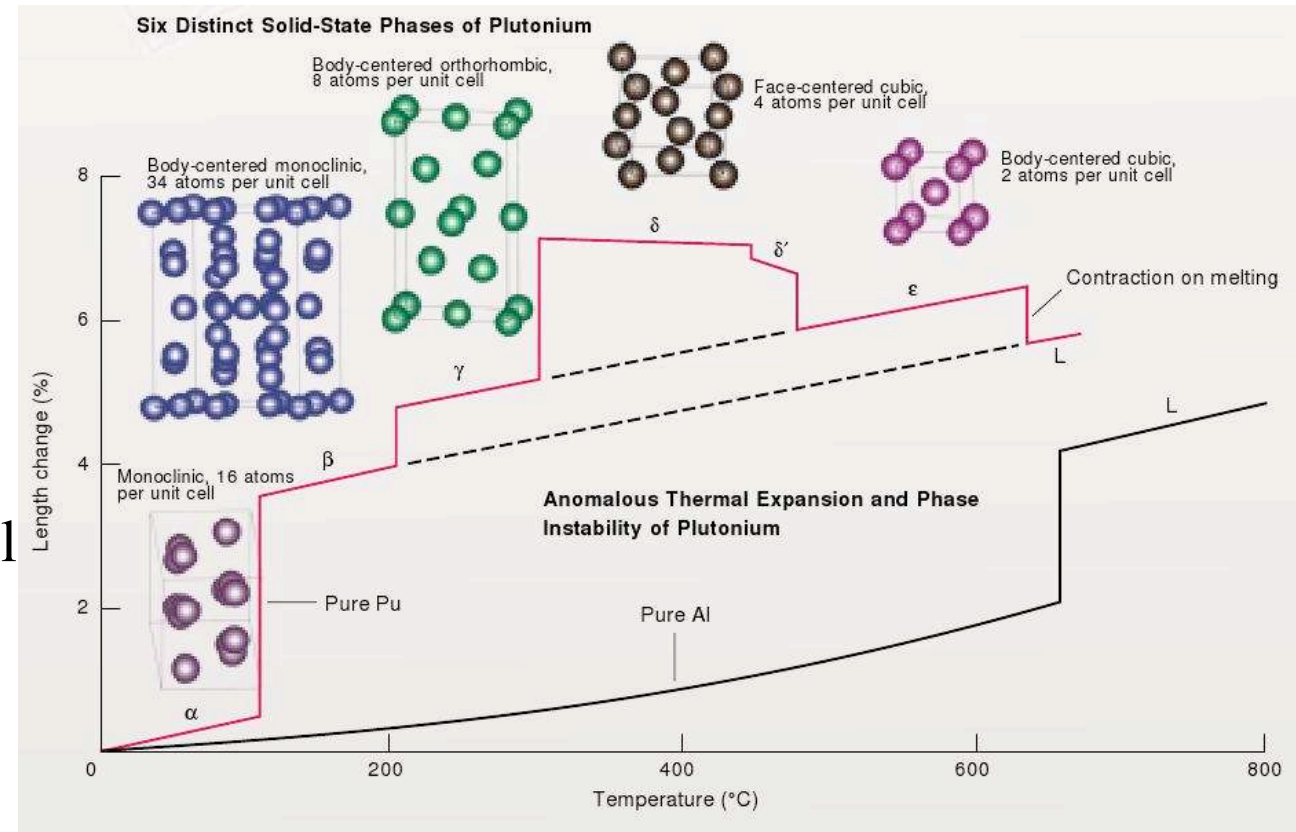
- Is Pu closer to atomic limit or band limit?
- Will one technique have convergence problems?
- Atomic limit performs significantly better.
 - Much better statistics per unit time.
- Why???
- Obvious how to truncate basis in atomic limit.
- Pu is actually closer to atomic-limit than band limit.

CTQMC atomic limit is preferential for Pu



Computational Objectives

- Approximation-free DMFT
- Ambient temperatures and below
- Non-trivial structures
 - More than one atom per unit cell
- Measure various observables



Computational Requirements

- QMC is rate limiting step of DFT+DMFT
 - Parallelizes nearly perfectly
 - More time \Rightarrow more progress
- Average run $\approx 20\text{k}$ CPU-hours
 - 1 atom/cell
 - scales linearly with # of atoms
 - $\approx 40\text{k}$ CPU-hours for $T = 240\text{K}$
- Self-consistency in density
 - factor of 10 increase?



Atlas

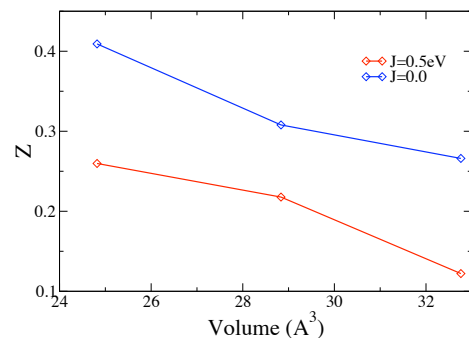
Results for δ Pu

- Performed first approximation-free DMFT

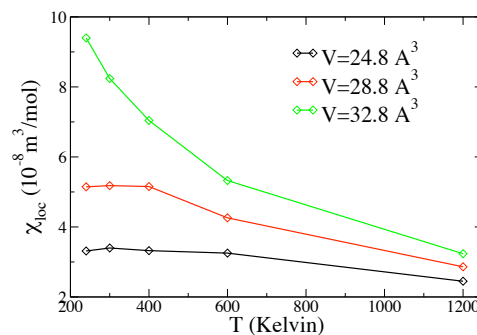
→ Including full exchange interaction

- Calculated various properties

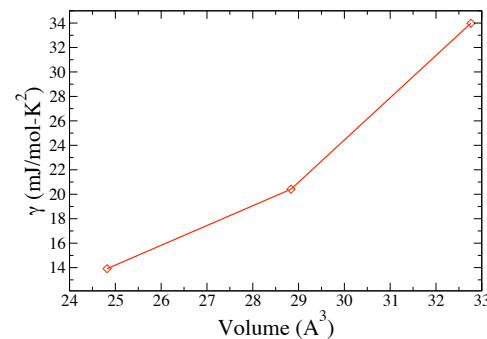
→ Quasiparticle weight



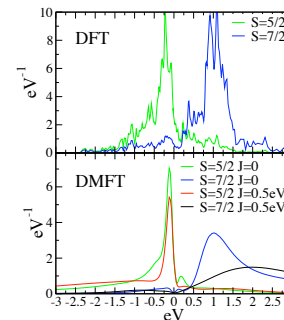
→ Magnetic susceptibility



→ Heat capacity

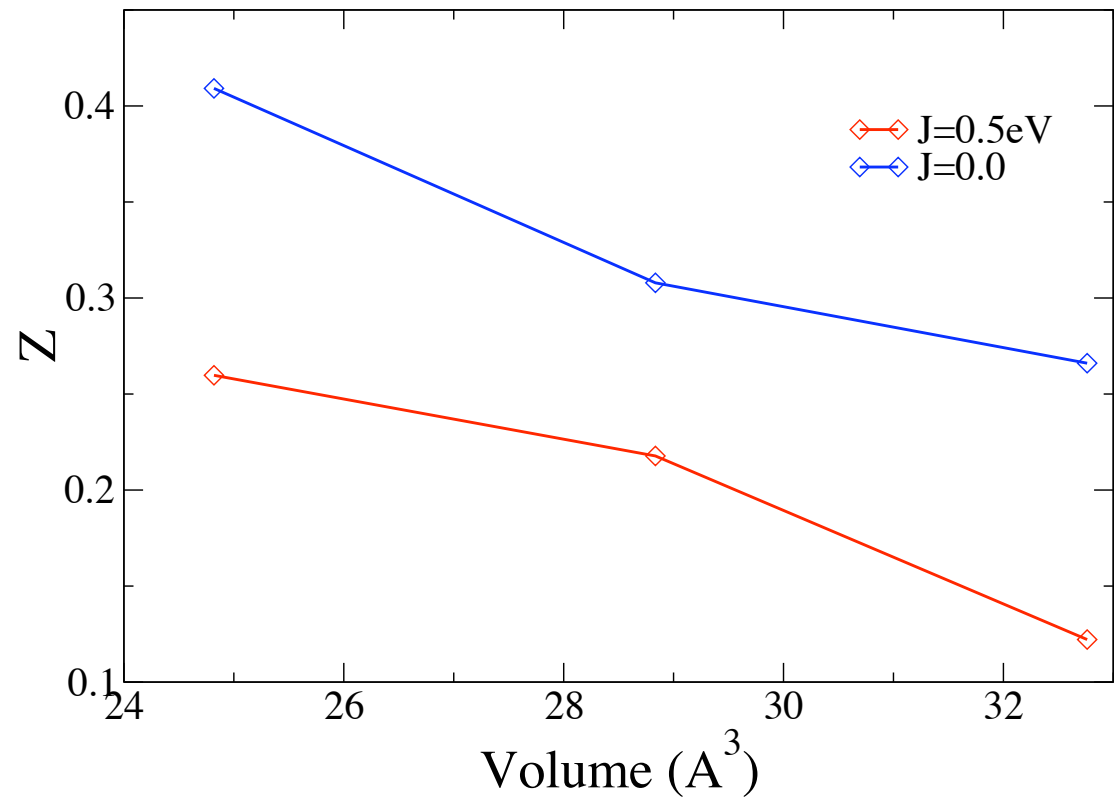


→ Spectra



Quasiparticle Weight $Z = \frac{1}{m^*}$

- Z is inverse of effective mass
 - $Z=1 \Rightarrow$ no correlations
 - $Z=0 \Rightarrow$ electrons localize
- Pu is closer to atomic limit
- Electrons become heavier as volume increases
- Including exchange substantially increases correlations.

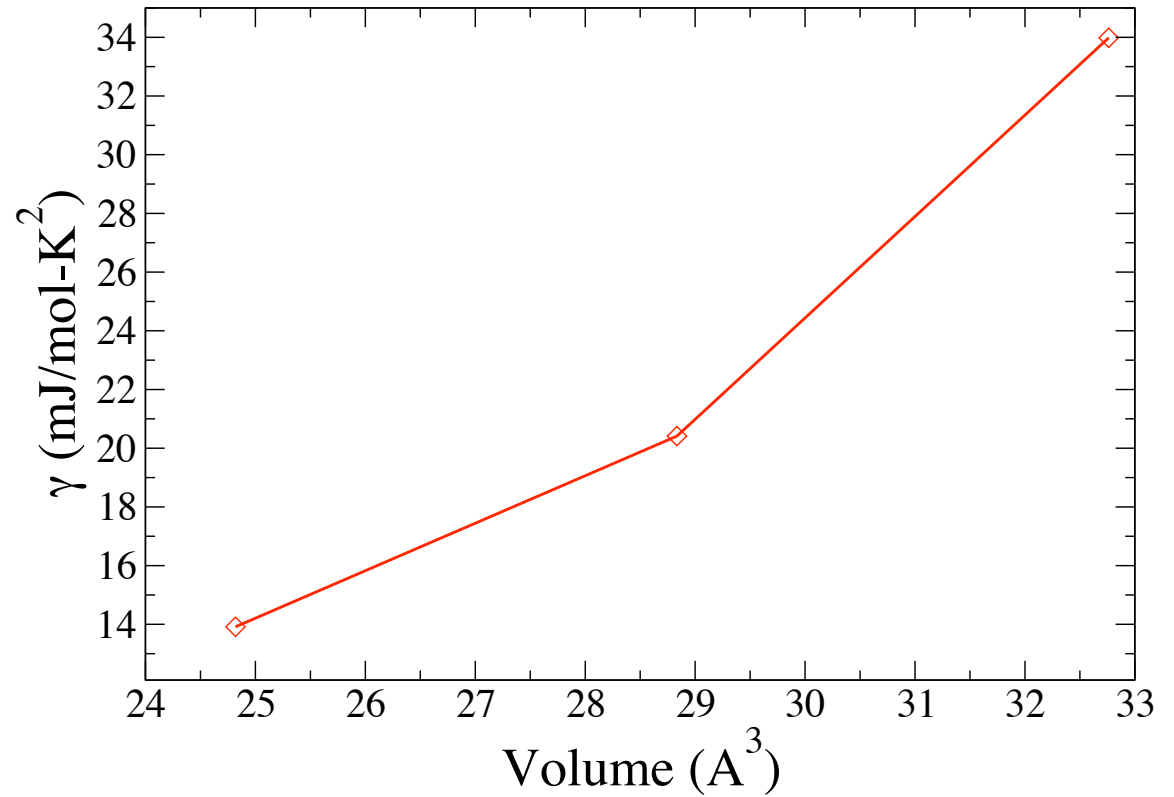


Pu is strongly correlated



Linear Coefficient of Heat Capacity

- Low temperature heat capacity sensitive to correlations.
- Strongly depends on volume.
- Experiments find $35 - 65 \frac{\text{mJ}}{\text{mol K}^2}$
 - Huge expt. variation
- Cause of discrepancy:
 - Density self-consistency
 - number of f electrons
 - Inaccurate experiments?

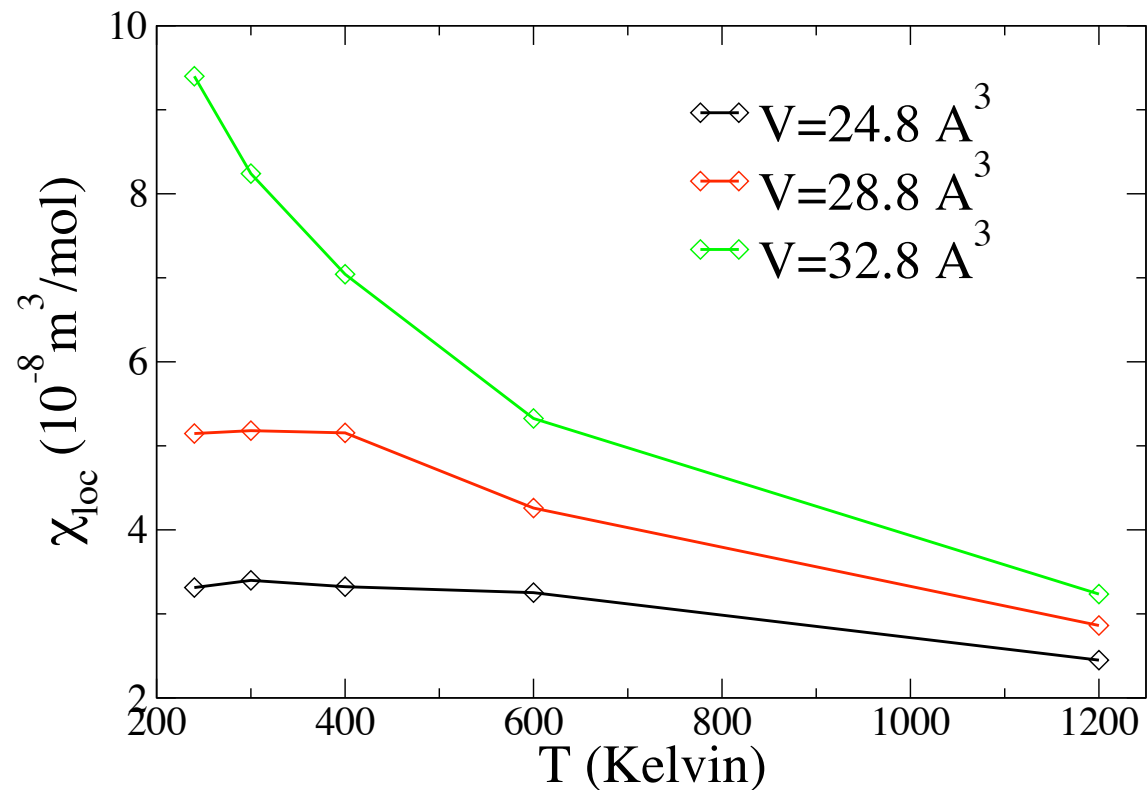


$$\gamma = \frac{2\pi k_B^2}{3} \sum_{\alpha} \frac{\rho_{\alpha}(0)}{Z_{\alpha}}$$



Magnetic Susceptibility χ

- First calculation of χ in Pu
- Pauli behavior \Rightarrow itinerant electrons
- Curie behavior \Rightarrow localized electrons
- Predict Pauli behavior for V_{eq}
 - \rightarrow Agrees with experiments
 - \rightarrow Explains lack of magnetism
- Expanded lattice agrees with PuH_2 .

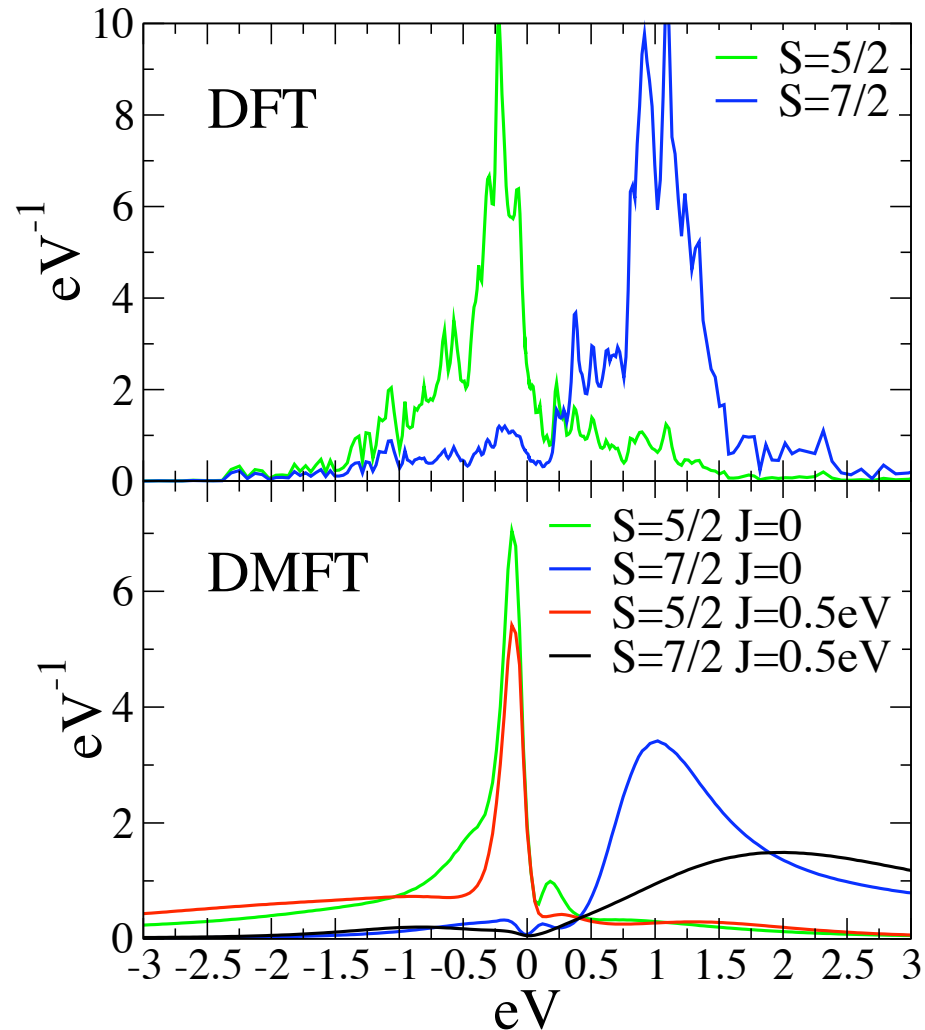


We have observed quantum decoherence in Pu



Photoemission spectra

- Density of electronic states versus energy.
- DMFT renormalizes DFT spectrum
 - Transfer of weight from Fermi energy to Hubbard bands
- Exchange causes further renormalization of spectrum.



Exchange has notable effect on spectra



Future Work

- Compute electronic properties of α Pu
 - 16 atoms per unit cell
- Compute total energy as a function of volume for δ Pu
 - Clearly elucidate volume collapse
- Compute negative thermal expansion
- Make corresponding predictions for Pu alloys.
 - Am, Ga, vacancies
 - Show why alloying equates to pressure



Computational and Scientific Impact

- First approximation-free exact DMFT calculation of Pu
 - *Impossible* without the power of Atlas
- Pu is shown to be a strongly correlated Fermi liquid
 - The moments in Pu are screened
 - DFT+DMFT is a pivotal tool for understanding strong correlated systems
- The use of massive parallel computation is critical to understanding the elusive properties of Pu.



Presentations / Publications

- Invited talk APS 2008 - UCRL-ABS-236899
 - *Dynamical mean-field theory calculations of materials properties using the continuous time quantum Monte-Carlo method*
- Invited talk Pu Futures 2008 - LLNL-ABS-401407
 - *Electronic coherence in δ Pu: A DMFT study*
- Contributed talk MRS 2008 - UCRL-ABS-236117
 - *Electronic properties of Pu via the dynamical-mean field theory*
- Contribute talk APS 2008 - UCRL-ABS-236889
 - *Effects of full Coulomb interactions on electronic structure of delta-Pu*
- Publication in preparation

